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# Effects of Various Shape Fillets On A 76/40 Double Delta Wing from Mach 0.18 to 0.7

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#### Nomenclature

b	Span (inches)
$\mathbf{BL}$	Buttline (inches)
c	Root chord (inches)
$C_{D}$	Drag coefficient
$C_{\mathrm{L}}$	Lift coefficient
$C_{m}$	Pitching moment coefficient at 73.27% root chord location
$C_{p,u}$	Upper surface pressure coefficient
MS	Model station (inches)
L/D	Lift-to-drag ratio
Re	Reynolds number based on root chord
x	Streamwise coordinate (inches)
у	Spanwise coordinate (inches)
α	Angle of attack (degrees)

## INTRODUCTION

The subsonic and transonic high angle-of-attack flow fields about advanced aircraft configurations (F/A-18, F-16, YF-22, YF-23, X-31, and JSF) are characterized by flow separation induced vortical flows, mutual interactions of multiple vortices, vortex breakdown, and interactions of the vortex flows with shock waves and control surfaces. The shedding of vortices from various locations on aircraft occur at the forebody, moderate to highly swept wing leading edges, strakes and leading edge extensions (LEX), and the junction of aerodynamic surfaces with each other and the fuselage. The shedding, interaction, and breakdown of these vortices are highly sensitive to both the aircraft's geometry and the flow conditions. In addition to producing the benefits of enhanced lift and maneuverability, the vortical flow also causes serious flight path departure and structural fatigue problems.<sup>1,2</sup> The ability to control the trajectory, strength, and breakdown processes of the vortices may not only alleviate the problems generated by this flow phenomena, but also enhance the existing controllability, maneuverability, and agility of the aircraft.

Because of the complexity and potential benefit from vortical flow, research has and is being conducted in this field to improve the understanding of the shedding, interaction, and breakdown processes of vortex flow.<sup>3-</sup>
This research has identified techniques for controlling vortical flows which provide additional control

power, complementing conventional control surfaces. These techniques enhance the maneuverability and broaden the envelope of aircraft operating conditions. The goal, however, is to efficiently exploit the potential benefits of vortex flow control.

The objective of this experimental investigation was to understand the impact of small geometry modifications (fillets) at the juncture of a 76° sweep leading edge extension (LEX) or strake and a 40° sweep cropped delta wing ("76/40 double delta wing"). A sharp leading edge model was used to depict a generic fighter strake and wing configuration. These fillets were developed to control the shedding, trajectory, and subsequent breakdown of vortices to enhance aircraft aerodynamic performance at subsonic and transonic speeds, and at elevated angles of attack. Longitudinal, lateral, and flow field characteristics of the double delta wing were investigated with different size linear, diamond, and parabolic shape fillets, Figures 1 and 2. The investigation was conducted in the 7x10 Ft. High Speed Tunnel at the National Aeronautics and Space Administration's (NASA) Langley Research Center (LaRC). The wind tunnel test was conducted from Mach 0.18 to 0.7; Reynolds numbers per foot from 1.2 x 10<sup>6</sup> to 3.6 x 10<sup>6</sup>; angles of attack from -4° to 42°; and zero sideslip. The aerodynamic characteristics of the baseline 76/40 double delta wing are discussed in The Baseline 76/40 Double Delta Wing section. The effects of the various fillets are discussed in Fillet Effects section.

## **BACKGROUND**

It is believed that by affecting the shedding process at thelized junction of the strake and wing, the trajectory, vortex-to-vortex interaction, and breakdown of the resulting vortex structure may be controlled resulting in enhanced lift and longitudinal control. Through asymmetric deployment of the fillets, additional lateral-directional control may be rea. The fillets are envisioned as being deployable or fixed at the junction of the strake (LEX) and wing. If deployable during maneuvering or landing approach, these fillets could enhance aerodynamic performance on demand.

## DOUBLE DELTA WING VORTEX FLOW FIELD

Several researchers have investigated the fundamental flow field of double delta wings with sharp leading edges. Cunningham and van Boer, <sup>7</sup> Brennenstuhl and Hummel, <sup>8</sup> Thompson, <sup>11</sup> Kern, <sup>14,15</sup> and Ekaterinaris et. al. <sup>16</sup> have conducted studies on double delta wings with various strake and wing leading edge sweeps. All the above studies had comparable results for similar leading edge sweeps. Regardless of leading edge sweep, a fundamental flow field exists across all double delta wings.

A double delta fundamental flow field is categorized as follows. At lower angles of attack, the strake and wing generate their own set of vortices that tend to follow the strake and wing leading edges. At the strake wing juncture, the strake vortex continues down stream no longer being fed by the vortex sheet from the strake. The wing generates a new stronger vortex (with increased circulation) than the strake vortex. As angle of attack increases, the wing vortex starts to move inboard away from the wing leading edge. With increasing angle of attack, the wing vortex entrains the strake vortex. At some angle of attack, the strake and wing vortices begin to coalesce. Depending on wing sweep, the wing vortex may breakdown (burst) before the strake and wing vortices begin coalesce. At higher angles of attack, both the wing and strake vortices burst. After the wing vortex burst point has traveled upstream with increasing angle of attack, it eventually reaches the strake/wing juncture. The strake vortex burst point moves further up stream, with increasing angle of attack, until it reaches the apex of the strake.

## VORTEX FLOW CONTROL HISTORY

The concept of flow control through vortex manipulation is not a new topic. H.R. Lawrence, in 1950, suggested the use of asymmetric edge shape effects to achieve roll control of low aspect ratio wings at high angles of attack.<sup>4</sup> In the 60's and 70's leading edge extensions (LEX) or strakes were developed and used in the development of F-5A/B, YF-17, F-16, and F/A-18. Several geometric modification concepts have been suggested by Lamar, <sup>18</sup> Rao, <sup>19, 20</sup> Campbell, <sup>20</sup> Bobbitt and Foughner, <sup>21</sup> and Malcolm. <sup>22</sup> Their concepts range from devices to enhance lift for landing and maneuvering to improving stability and control (i.e. hinged

strakes, vortex flap, apex flap, leading edge vortex generator, etc.). References 14 and 15 provide a comprehensive history of vortex flow control.

A computational investigation on the effect of the fillets presented in this paper was conducted by Kern. Straka, while investigating the effects of wing planform on vortex break down, developed concave fillets similar to the parabolic fillets developed by the Naval Air Systems Command and Office of Naval Research "Vortex Flow Control" program. The Naval Postgraduate School conducted a water tunnel test to investigate the highly interactive vortical flow field generated by the fillets presented in this paper. A conducted a water tunnel test to investigate the highly interactive vortical flow field generated by the fillets presented in this paper.

## EXPERIMENTAL INVESTIGATION

The wind tunnel test was conducted with a model depicted a generic fighter or high speed civil transport. The model, presented in Figure 1, is comprised of a flat plate cropped double delta wing with a 76° strake (LEX) and a 40° wing. The strake and wing leading edges are beveled 20° and sharp. The chord and wing span are 16 and 16.326 inches, respectively. The strake/wing juncture is located at MS 8.933 and BL 2.227; the cropped portion of the wing begins at MS 13.46 and BL 8.163. The model is 0.375 inches thick.

To control the vortex flow, three fillet shapes were attached at the strake and wing juncture, of the baseline configuration. The parabolic fillets removed the discontinuity of the junction, while the linear and diamond-shaped fillets added additional discontinuities, from which additional vortices were expected to originate. Three fillet sizes based on total fillet area (left and right) to model reference area were tested; 0.5%, 1%, and 2.5% for linear and diamond, and 0.5%, and 1% for parabolic. The various 1% fillets are illustrated in Figure 2. Both symmetrical and asymmetric fillets were tested. However, only symmetrical data will be discussed.

The model's six component forces and moments were measured with a bottom-mounted sting balance. The moment reference center (MRC) of the double delta wing was located at 73.27% of the root chord. This MRC coincides with the MRC used in References 7 and 15 on a similar double delta wing planform. The balance was covered with a fairing. Balance pitch attitude was measured with an accelerometer mounted near the sting's pivot point of the model support system. Roll was determined with a potentiometer installed in the model support system. All angle measurements were corrected for balance and sting deflections. Forty-seven surface static pressure orifices (taps) were used during the test. They were distributed in three spanwise rows on the left-hand side of the upper surface of the wing. The spanwise rows were located at x/c = 0.25, 0.75, and 0.9. A forty-eight ESP module house in the fairing forward of the balance was used to measure orifice pressures.

## PRESSURE SENSITIVE PAINT

To supplement surface pressure orifice data, a platinum octaethylporphyrin-based luminescent pressure sensitive paint (PSP) developed by NASA Ames Research Center and The University of Washington was applied to the upper surface of the model. The painted model was illuminated by six ultraviolet (UV) lamps and photographed by an 8-bit CCD camera mounted outside the test section ceiling. The "raw" wind-off and wind-on PSP images were registered and resectioned into a planform view by software developed at NASA Ames. The paint provided the pressure distribution over the entire upper surface of the wing. All PSP images show black bands where no paint was applied. To determine paint thickness and texture effects on surface static pressure, measurements from the pressure orifices were taken with and without PSP. PSP and pressure tap data were taken simultaneously. PSP surface finish effects on pneumatic data were very small. All upper surface static pressure orifice values presented in this paper were taken without pressure sensitive paint on the model.

## FLOW VISUALIZATION

In addition to force, moment, and pressure measurements, off-body flow visualization was conducted with a vapor screen technique.<sup>27</sup> The ambient relative humidity in the test section created condensation within the vortical flow of the model. The vortex cross sections were illuminated with a laser light sheet revealing bright vortices with a dark background. Light sheet position was varied remotely to provide coverage of the model.

The light sheet optics were located above the test section ceiling. The off body flow field features were recorded by video cameras mounted on the sting and outside the test section.

# WIND TUNNEL FACILITY AND TEST CONDITIONS

The model force, moment, surface static pressures, and surface pressure sensitive paint measurements and laser vapor screen results were obtained in the 7x10 Ft. High Speed Tunnel (HST) at NASA Langley Research Center in Hampton, Virginia. The 7x10 Ft. HST was a continuous-flow, closed circuit, and atmospheric facility capable of operating to Mach 0.8. A complete description of the HST is provided in Reference 28.

The test was conducted at free stream Mach numbers 0.18, 0.3, 0.5, and 0.7 and angles of attack from -4.0° to 42.0°. Angle of attack was limited to approximately 28° and 16° at Mach 0.5 and 0.7, respectively, due to balance limits. No sideslip conditions were investigated. The tunnel stagnation pressure was set by atmospheric conditions. The Reynolds numbers per foot were approximately 1.2 x10<sup>6</sup>, 2.0 x10<sup>6</sup>, 3.0 x10<sup>6</sup>, and 3.6 x10<sup>6</sup>. No base or balance cavity pressures were measured. Solid blockage and wall constraints corrections were applied. The strake, wing, balance fairing, and fillets were tested with free transition.

# **DISCUSSION OF RESULTS**

# BASELINE MODEL LONGITUDINAL AND FLOW FIELD CHARACTERISTICS

# Baseline Longitudinal Characteristics

The baseline double delta wing longitudinal characteristics at M=0.18, 0.3, 0.5, and 0.7 are presented in Figure 3. The lift, drag, and pitching moment coefficient for M=0.18 and 0.3 are characterized by two points where there is a slope change (break). The first slope break is associated with the wing vortex bursting at or near the wing apex. The second slope change (stall) is associated with the strake vortex bursting close the strake's apex. At M=0.18 and 0.3, the first slope change occurs at 14° rather than at 16° for M=0.5. The M=0.5 lift curve collapses onto the M=0.18 and 0.3 curves at 24°. The M=0.7 data does not change slope abruptly between -4° to 16° angle of attack. A lift slope increase associated with compressibility is observed at Mach 0.5 and 0.7. There is a slight reduction in induced drag for a given lift coefficient with increasing Mach number. Pitching moment stability also increases with increasing Mach. A pitch-up associated with the wing vortex bursting occurs at the first lift break. Pitching moment decreases along the same curve after the strake vortex bursts. Table 1 summarizes the longitudinal coefficient characteristics. Baseline repeatability is shown in Figure 4. Data repeated below 16° angle of attack. Above 16° angle of attack the data repeated with varying level of success. It is not know if Reynolds number variations or slight inconsistencies in model changes cause the scatter at higher angles of attack.

	Values at First Break Point				Values at Second Break Point (Stall)			
Mach	Alpha	$C_{L}$	$C_{D}$	$C_{\rm m}$	Alpha	C <sub>Lmax</sub>	$C_{D}$	$C_{m}$
0.18	14°	0.625	0.2	0.075	36°	0.87	0.74	0.137
0.30	14°	0.6	0.2	0.075	36°	0.82	0.7	0.125
0.50	16°	0.7	0.25	0.076	N/A	N/A	N/A	N/A
0.70	N/A	N/A	N/A	N/A	N/A	N/A	N/A	N/A

Table 1- Longitudinal Data Slope Change (Break) Points

## Baseline Flow Field Characteristics

Figures 5 and 6 present corresponding surface pressure distribution plots and PSP images at Mach 0.5 and 0.7. Vapor screen images supplement the pressure data. At Mach 0.5 and 10° angle of attack, the strake and wing vortices coexist as individual vortices with no mixing. The wing and strake vortices are mixing/intertwining near the trailing edge at 14° angle of attack. At 14° angle of attack two distinct footprints are seen at x/c of 0.75. However, a coalescing footprint is present at a x/c of 0.9 indicating the wing and strake vortices coalesce into a rolled-up vortex. As angle of attack is increased, the strake and wing vortex pressure footprints coalesce into a single footprint and burst further upstream. At 20° angle of attack, as single vortex

pressure footprint is present at x/c of 0.75 and 0.9. The PSP data shows the wing vortex generates a large amount of suction near the wing/strake juncture. Mach 0.7 data Figure 6, shows two distinct strake and wing vortices up to 16° angle of attack. The pressure footprints from the wing vortices appear diffused, indicating the wing vortices may be bursting. The Mach 0.7 PSP images show dark color striations on the strake vortex pressure footprint. The striation pattern has the wrong orientation to be the high and low suction peaks associated with the shear layer instability roll up shown in Reference 12, Figure 16 and in Reference 29, Figures 3.22 and 3.23. As depicted in References 5 through 16 the near surface flow under a vortex is downstream and spanwise. The striations are perpendicular to the near surface flow under a vortex, supporting the hypnosis the striations are caused by localized shock waves and subsequent pressure recovery. Along with the striations the vapor screen images show a ring of condensation around the primary condensations pattern of the strake vortex. It is not known if this is a connecting vortex sheet between the strake and wing vortex sheets.

# FILLET EFFECTS ON LONGITUDINAL AND FLOW FIELD CHARACTERISTICS

Fillet Geometry Effects

The linear, diamond and parabolic fillets increase lift and reduce drag for a given  $C_L$ . The diamond and linear fillets have a small effect on pitching moment. The parabolic fillets delay pitch-up. However, the parabolic fillets pitch-up is larger than the other fillets or the baseline configuration. The linear fillet increases  $C_L$  from 2.9% to 28.6%. The diamond fillet increases  $C_L$  from 2.9% to 17.9%. The parabolic fillet increases  $C_L$  from 20% to 39.3% depending on Mach and fillet size.

- M = 0.18

A comparison of the 0.5% linear, diamond, and parabolic fillets to the baseline configuration at Mach 0.18 is provided in Figure 7. The parabolic fillet generates a higher lift coefficient and lower drag (for a given  $C_L$ ) after 12° as compared to the baseline. The parabolic fillet pitching moment also is more stable than other fillets pitching moment. In turn, the parabolic fillet generates a higher pitch-up and larger lift slope change after the vortex's burst location moves over the wing. The diamond fillet lift coefficient is higher than the linear fillet lift coefficient from 16° to 24°. The linear fillet has a higher lift coefficient than the diamond fillet from 26° to 36°. There is a similar crossover in drag between the diamond and linear fillets. The diamond fillet generates a smoother lift curve than the linear or parabolic fillet. Both the linear and diamond fillets generate similar pitching moment curves.

- M = 0.30
- The 0.5% linear, diamond, and parabolic fillets are compared to the baseline model at Mach 0.30 in Figure 8. The parabolic fillet lift is higher than the diamond or linear fillets. However,  $C_{Lmax}$  decreases from 1.75 to 0.99 with increasing Mach. Unlike the Mach 0.18 results, the linear fillet generates an equal or higher lift coefficient than the diamond fillet. The linear fillet does not change lift slope until 18° ( $C_L$ = 0.8) when compared to 14° at Mach 0.18. However, the lift slope change is abrupt. The linear fillet drag coefficient tracks the parabolic fillet up to a  $C_L$  of 0.8. The linear fillet pitching moment coefficient matches the parabolic fillets pitching moment up to  $C_L$ = 0.8. The Mach 0.3 parabolic fillet pitching moment is less stable than at Mach 0.18.
- M = 0.50

The parabolic fillet generates the most lift followed by the linear fillet, diamond fillet, and baseline configurations, as shown in Figure 9. Both the parabolic and linear fillets have a sharp lift loss after 18°. The diamond fillet has the same lift curve shape as the baseline configuration up to 20° with 2.9% greater lift at 18°. At 21° angle of attack the baseline model lift curve intersects the diamond fillet's lift curve. The parabolic and linear fillets have drag coefficient curves with sharp breaks similar to the lift curve. Data was not obtained beyond 24° due to balance limits.

- M = 0.70

The Mach 0.70 data, Figure 10, is limited to 16° due to balance limits. The parabolic fillet generates a higher lift slope, lower drag and pitching moment than the other fillets and the baseline configuration. Below 15° angle of attack the linear fillet generates the same lift, drag and pitching moment than the baseline

configuration. At 16° (balance limit) the linear fillet deviates from the baseline configuration generating a more lift than the baseline configuration. The parabolic fillet closely follows the baseline configuration

## - Repeatability

The diamond fillet repeatability is similar to the baseline configuration. The linear fillet repeatability is presented in Figure 11. Two of the three runs have good agreement; however, the second run shows a large discrepancy. The second repeat run has a lower lift and higher pitching moment from 12° to 18° angle of attack. It is not known if two by-stable vortex structures exist or if run-to-run model change inconsistencies cause the discrepancies in the lift curves.

#### Inverted Model

Unexpected results were obtained when the model was inverted. As shown in Figure 12, the linear and diamond fillets did not increase lift a substantial amount when compared to the upright data. The lift coefficient of the inverted parabolic fillet is also lower than the upright model lift. It is not known if the decrease between the upright and inverted model is caused by balance fairing acting as a fuselage that decouples the port and starboard vortices, or whether the change in leading edge bevel shape affects the vortex sheet roll-up characteristics.

# Fillet Size Effects On Longitudinal Characteristics

The effect of fillet shape and size (0.5%, 1%, and 2.5%) on longitudinal characteristics are presented at Mach 0.18, 0.3, and 0.5. Effects of fillet shape and size on flow field characteristics are discussed after longitudinal characteristics.

## - Linear Fillet

The effect of the 0.5%, 1%, and 2.5% linear fillets on longitudinal force and moment characteristics at Mach 0.18, 0.3, and 0.5 are shown in Figures 13 through 15. Figure 13 shows a trend of increasing lift with increasing fillet size at Mach 0.18. Increasing fillet size does not have a large influence on lift at or above 30° angle of attack. Drag coefficient increases with increasing fillet size for a given angle of attack. However, it decreases for a given C<sub>L</sub>. The 0.5% and 1% linear fillets increase pitching moment instability slightly. There was a more pronounced instability increase with the 2.5% linear fillet. There is no clear correlation with increasing Mach. One fillet size will generate a higher lift coefficient within an alpha range at one Mach but behave differently at another Mach. As depicted in Figures 13 through 15 as Mach increases, fillet size effects diminish.

## - Diamond Fillet

The effect of the 0.5%, 1%, and 2.5% diamond fillets on longitudinal force and moment characteristics at Mach 0.18, 0.3, and 0.5 are shown in Figures 16 through 18. The only clear effect of diamond fillet size is at Mach 0.3 between 14° and 22° angle of attack. The 2.5% diamond fillet generates a highest lift coefficient from 14° to 22°. The 1% diamond fillet has a higher  $C_L$  than the 0.5% diamond fillet from 16° to 22°. Only the 2.5% fillet shows an appreciable increase of pitching moment instability.

#### - Parabolic Fillet

Longitudinal data is presented in Figures 19 through 21 at Mach 0.18, 0.3, and 0.5 respectively. Increasing fillet size from 0.5% to 1% creates discontinuities in the  $C_L$ ,  $C_D$ , and  $C_M$  curves. The 0.5% and 1% fillets generally produce the same force and moment coefficients up to the first discontinuity. The angle of attack of first  $C_L$  discontinuity decreases with increasing fillet size. The parabolic fillet is prone to large roll dynamics caused by unstable (oscillatory) left/right vortex bursting. At some angles of attack, the left and right vortices burst in a fixed symmetric or asymmetric pattern.

## Fillet Flow Field Characteristics

Flow field characteristics will be discussed by presenting upper surface pressure distribution plots from pneumatic data and PSP images. Vapor screen images will also be used in conjunction with the pressure data. The pressure data is presented at 8°, 12°, 16°, 20°, 22°, and 24° angle of attack. The pressure data will be limited to the 1% fillets and Mach 0.5.

# - Linear Fillet

The linear fillet flow field is dominated by the strake, fillet, and wing vortices. The strake vortex is immediately entrained by the fillet vortex causing the pneumatic and PSP data, illustrated in Figure 22, to show the fillet and wing vortices as a wide footprint with two suction peaks. At 16° angle of attack, the vortices mutually interact with which other resulting in a single footprint. The vapor screen images show the strake, fillet, and vortices roll-up into a single structure with increasing angle of attack. The fillet vortex is entrained under the wing vortex while the fillet vortex entrains the strake vortex. The strake and fillet vortex cores rotate around the wing vortex core.

## - Diamond Fillet

Four pairs of vortices dominate the diamond fillet flow field: one from the strake, two from the fillet, and one from the wing. The vortex structure at various model stations is presented in Figure 23. The two vortices, emanating from the fillet, rapidly roll-up into a single vortex. Since the fillet vortices roll-up into a single vortex, the upper surface pressure plots only show three distinct footprints at low angles of attack. As shown in the pressure tap data and PSP images, the strake, fillet, and wing vortices start to show a single pressure footprint with two suction peaks at higher angles of attack. The PSP images show the left wing vortex bursting at 22° angle of attack and the strake vortex pressure footprints becoming asymmetric.

## - Parabolic Fillet

The pneumatic, PSP, and vapor screen data in Figure 24 show the majority of the vorticity rolls up into a single pair of vortices (left and right). This is also depicted by a single pressure peak. The vapor screen images show one vortex pair for the majority of vortex trajectory. The synergistic effect of concentrating the vorticity into a single vortex creates large suction peaks. The combined strake/wing vortex ceases to follow the wing leading edge and "tares off." After the vortex tares, a new vortex is formed along the wing leading edge. Asymmetric vortex bursting is clearly shown in the PSP data at 22 and 24 angles of attack.

#### SUMMARY AND CONCLUSIONS

A 76°/40° sweep double delta wing was tested in the NASA Langley 7x10 ft HST from Mach 0.18 to 0.7 and from -4° to 42° angle of attack. The double delta wing strake/wing juncture was modified with three fillet shapes: linear, diamond, and parabolic. The linear and diamond fillet size varied 0.5%, 1%, and 2.5% (left and right fillet area/ wing planform area). The parabolic fillet size varied 0.5% and 1% in size.

Baseline configuration lift coefficient increased while induced drag coefficient decreased with increasing Mach. Since the wind tunnel Reynolds number varied with atmospheric conditions and Mach, Reynolds number effects can not be separated from Mach effects. Pitching moment coefficient did not vary below Mach 0.5 and a  $C_L$  of 0.6. At Mach 0.7 pitching moment was more stable. Above a  $C_L$  of 0.6 there is no clear Mach trend. The lift curves can be characterized by three distinct regions: 1) an angle of attack range where lift slope gradually increases, approximately from  $-4^{\circ}$  to  $12^{\circ}$  angle of attack. In this angle of attack range, the wing vortex does not entrain and mix with the strake vortex. 2) From approximately  $12^{\circ}$  to  $32^{\circ}$  angle of attack, the lift curve can be characterized with a break and reduction in lift slope. In this angle of attack range, the wing vortex entrains the strake vortex. The vortex burst point of the combined vortices moves upstream with increasing angle of attack. At higher angles of attack, the strake vortex bursts over the strake while the wing vortex bursts at or near the strake/wing juncture. 3) A post stall region, after  $36^{\circ}$  angle of attack, where lift slope is negative. The strake vortex burst point is near the strake apex. Increasing Mach delays the entertainment of the strake vortex by the wing vortex.

Baseline configuration PSP images at Mach 0.7 clearly show dark striations on the strake vortex footprint. The striation color indicates regions of local high pressure. The striation pattern has the wrong orientation to be the high and low suction peaks associated with the shear layer instability roll up shown in Reference 12, Figure 16 and in Reference 29, Figures 3.22 and 3.23. As depicted in References 5 through 16 the near surface flow under a vortex is downstream and spanwise. The striations are perpendicular to the near surface flow under a vortex, supporting the hypnosis that the striations are caused by localized shock waves and subsequent

pressure recovery. The striations are not seen in the wing vortex. At present it is not known what causes the striations.

All the fillets increase lift from approximately 8° to 32° angle of attack. The parabolic fillet generates the most lift, while the linear and diamond fillet lift increments vary at different angles of attack. Induced drag decreases with increasing fillet size. Pitching moment does not significantly vary with the linear and diamond fillets. The parabolic fillet increases pitch stability at low angles of attack and pitch-up at elevated angles of attack. Increasing fillet size typically increases lift and pitch stability, and reduces induced drag. However, increasing parabolic fillet size also increases the number of lift curve breaks and lowers the angles of attack where they occur.

The fillet effects decrease with increasing Mach. It is hypothesized this is caused by a delay of the strake vortex entertainment by the wing vortex. A similar delay between strake and wing vortex interactions is observed in the baseline configuration.

The fillets show potential of improving aircraft aerodynamic performance. However, further research needs to be conducted on the Reynolds number, center body (fuselage), and leading-edge shape effects on fillet effectiveness. The authors Gonzalez and Erickson conducted a subsequent test (unpublished) to evaluate the impact Reynolds number has on fillet effectiveness.

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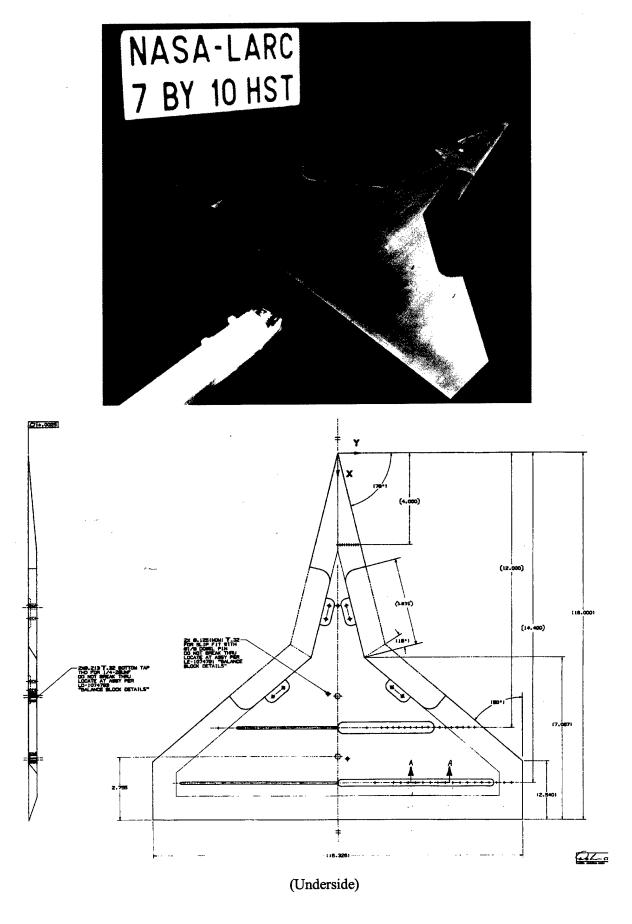


Figure 1- Model in 7 by 10 ft HST and Model Details

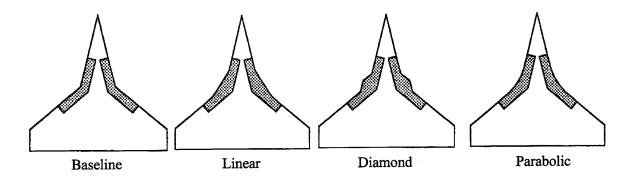


Figure 2- Baseline model and fillet configurations

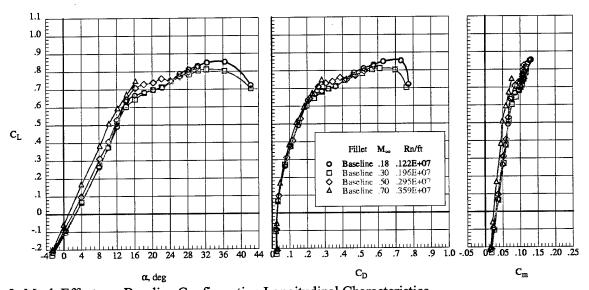


Figure 3- Mach Effects on Baseline Configuration Longitudinal Characteristics

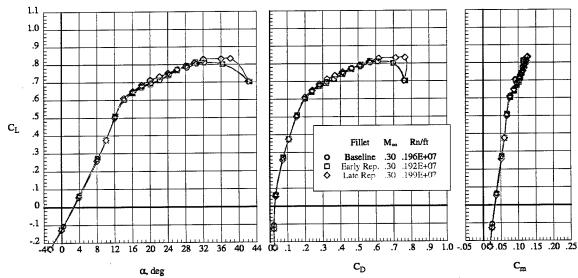


Figure 4- Repeatability of Baseline Configuration Longitudinal Characteristics, Mach 0.3

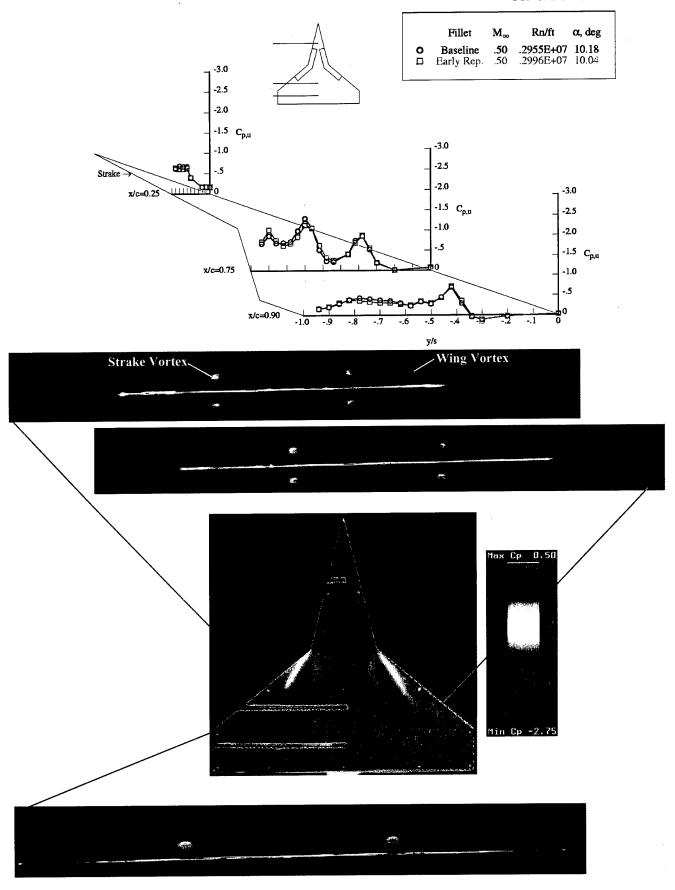


Figure 5- Baseline Configuration Flow Field and Pressure Distribution, Mach 0.5, AOA 10°

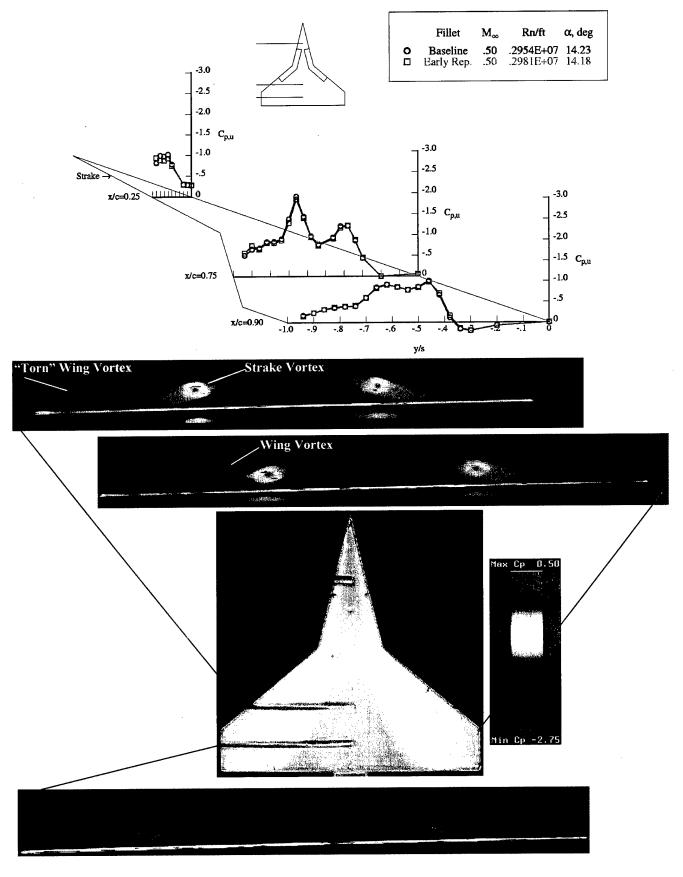


Figure 5 (cont.)- Baseline Configuration Flow Field and Pressure Distribution, Mach 0.5, AOA 14°

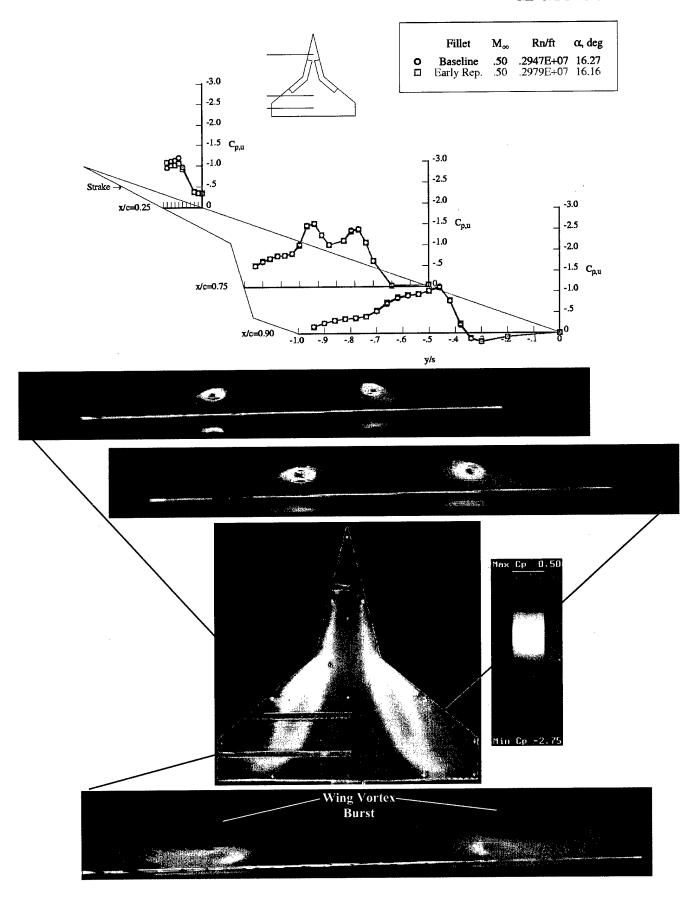


Figure 5 (cont.)- Baseline Configuration Flow Field and Pressure Distribution, Mach 0.5, AOA 16°

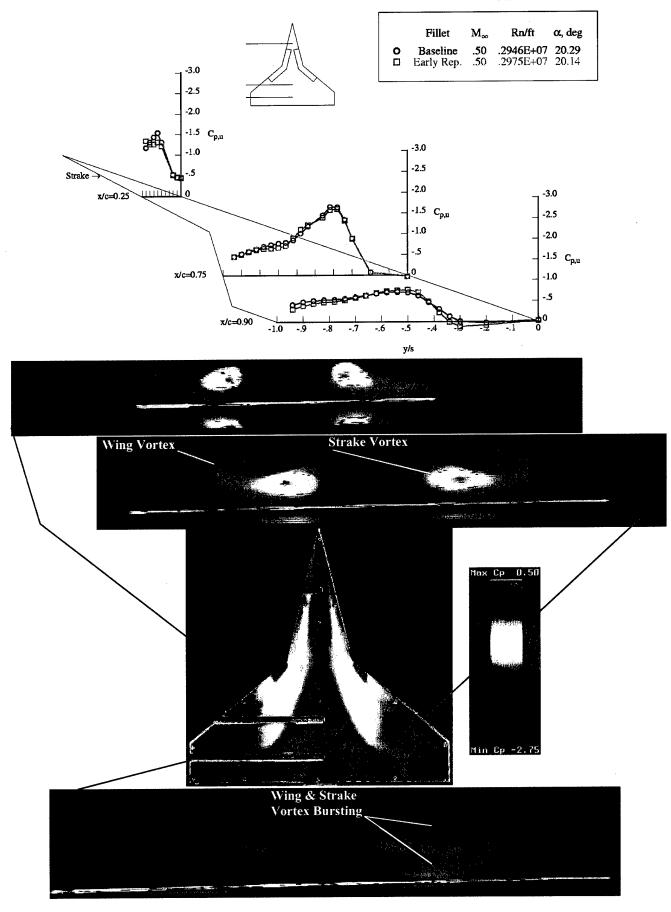


Figure 5 (cont.)- Baseline Configuration Flow Field and Pressure Distribution, Mach 0.5, AOA 20°

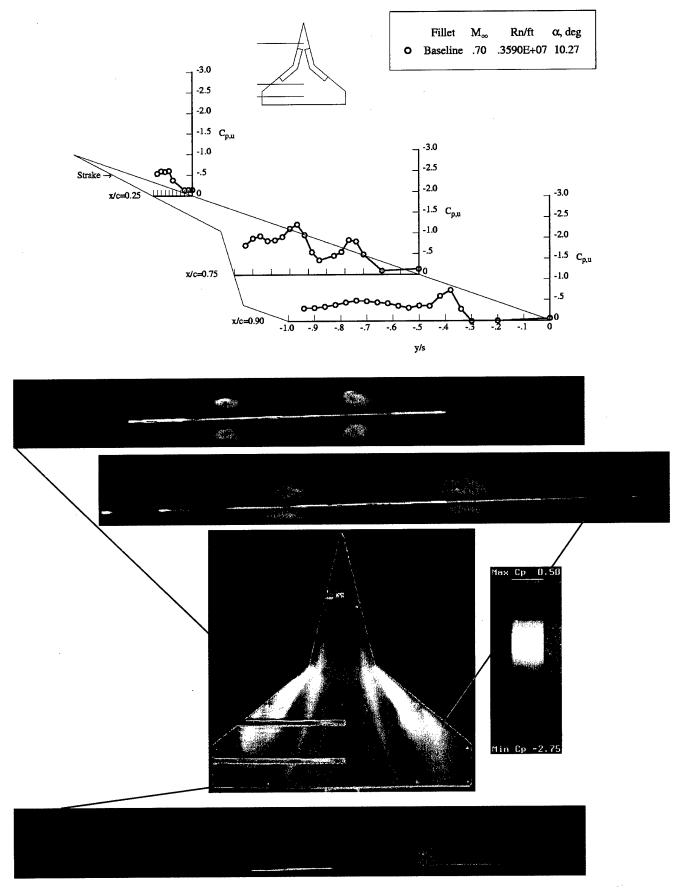


Figure 6- Baseline Configuration Flow Field and Pressure Distribution, Mach 0.7, AOA 10°

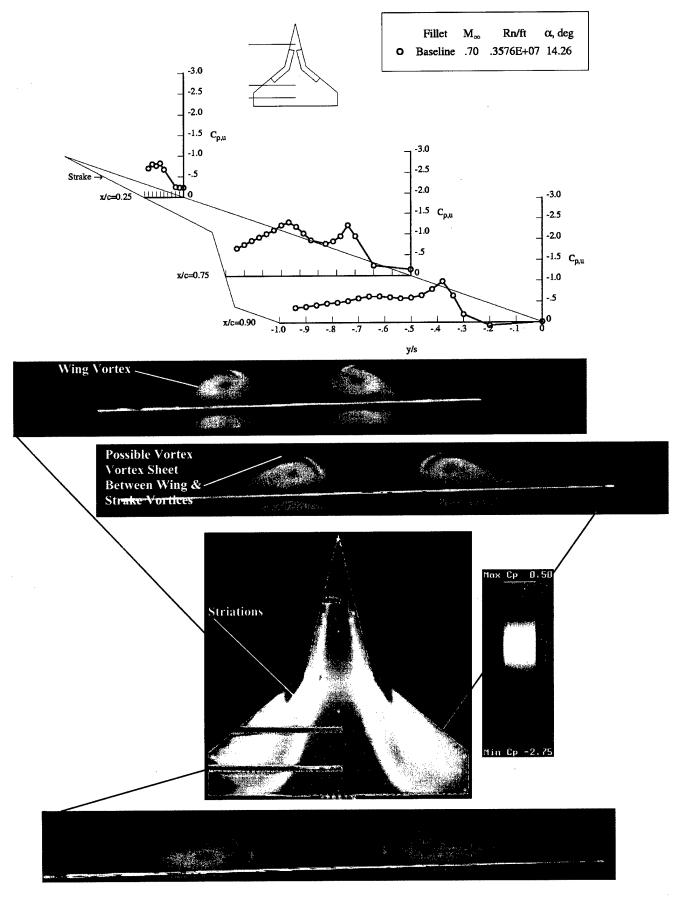


Figure 6 (cont.)- Baseline Configuration Flow Field and Pressure Distribution, Mach 0.7, AOA 14°

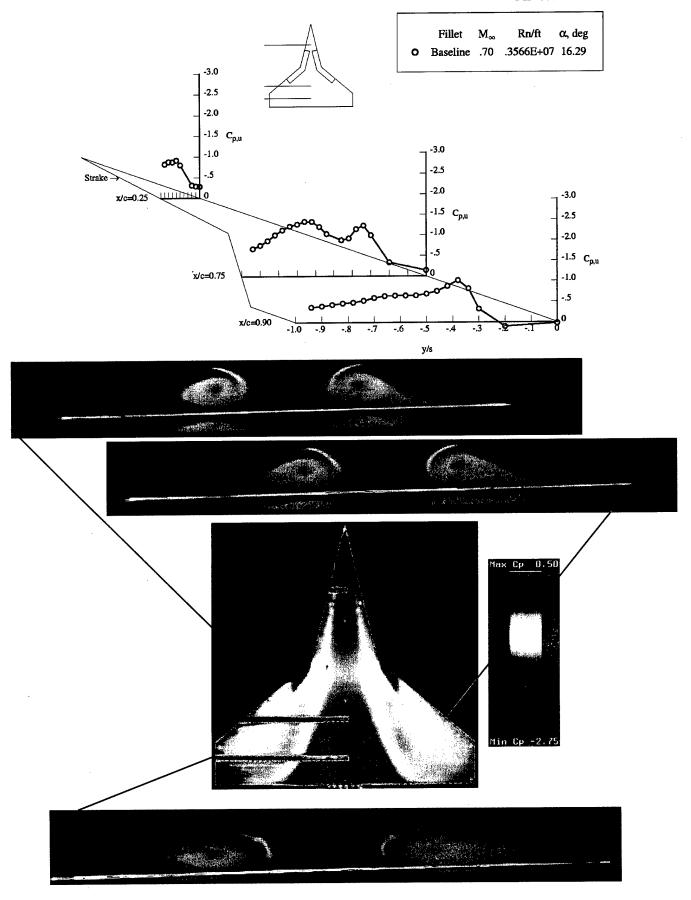


Figure 6 (cont.)- Baseline Configuration Flow Field and Pressure Distribution, Mach 0.7, AOA  $16^\circ$ 

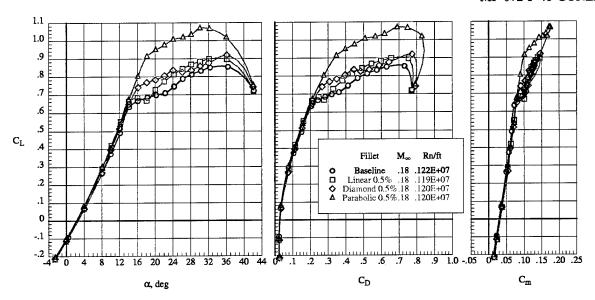


Figure 7-0.5% Fillet Effect on Longitudinal Characteristics, Mach 0.18

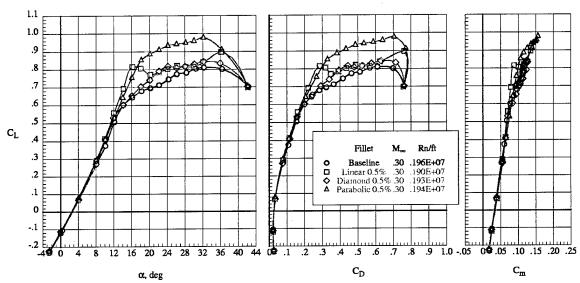


Figure 8-0.5% Fillet Effect on Longitudinal Characteristics, Mach 0.3

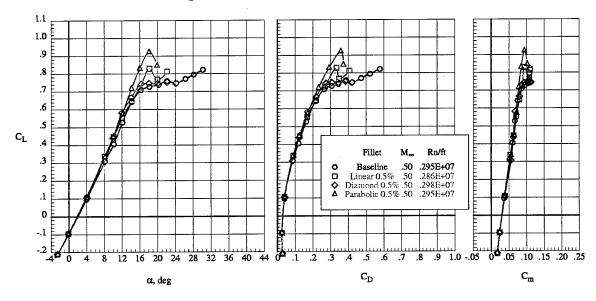


Figure 9- 0.5% Fillet Effect on Longitudinal Characteristics, Mach 0.5

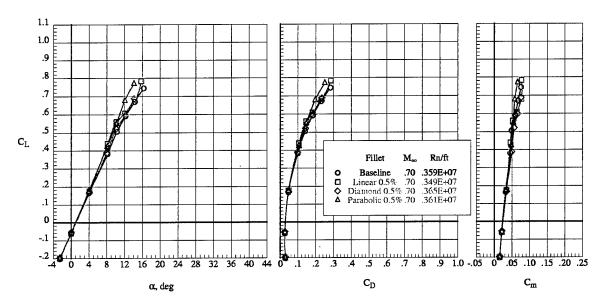


Figure 10-0.5% Fillet Effect on Longitudinal Characteristics, Mach 0.7

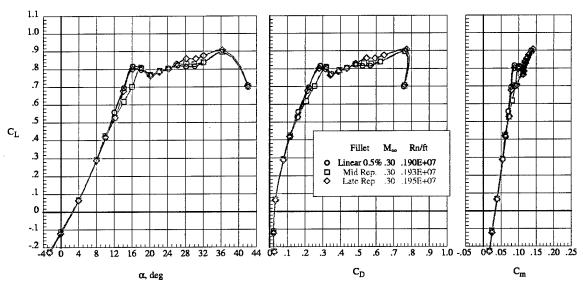


Figure 11-1% Linear Fillet Repeatability, Mach 0.3

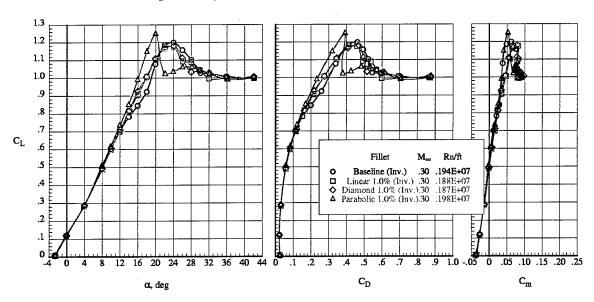


Figure 12- Fillet Effects On Inverted Model, Mach 0.3

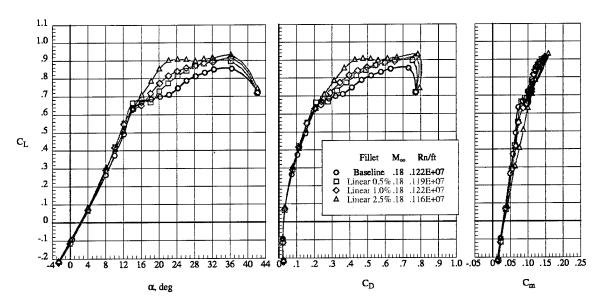


Figure 13- Linear Fillet Size Effect on Longitudinal Characteristics, Mach 0.18

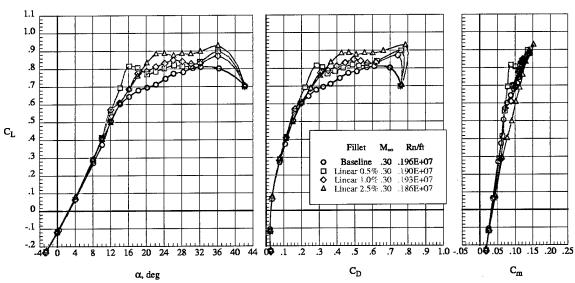


Figure 14- Linear Fillet Size Effect on Longitudinal Characteristics, Mach 0.3

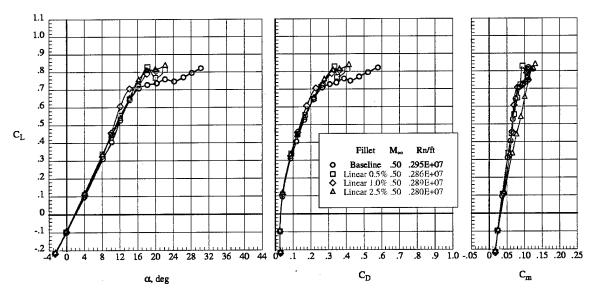


Figure 15- Linear Fillet Size Effect on Longitudinal Characteristics, Mach 0.5

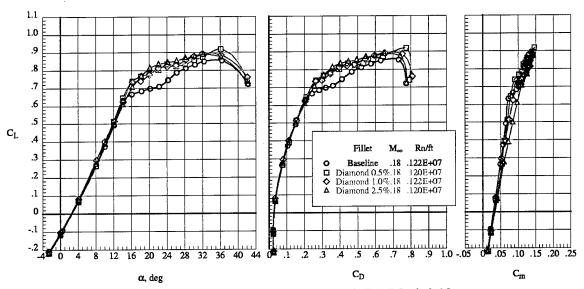


Figure 16- Diamond Fillet Size Effect on Longitudinal Characteristics, Mach 0.18

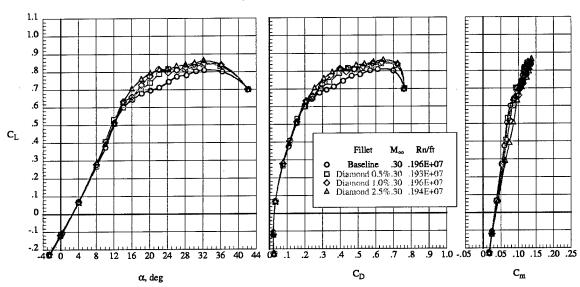


Figure 17- Diamond Fillet Size Effect on Longitudinal Characteristics, Mach 0.3

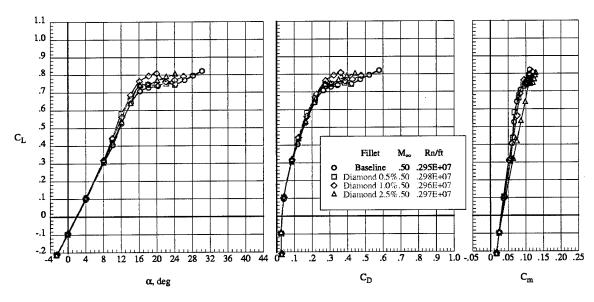


Figure 18- Diamond Fillet Size Effect on Longitudinal Characteristics, Mach 0.5

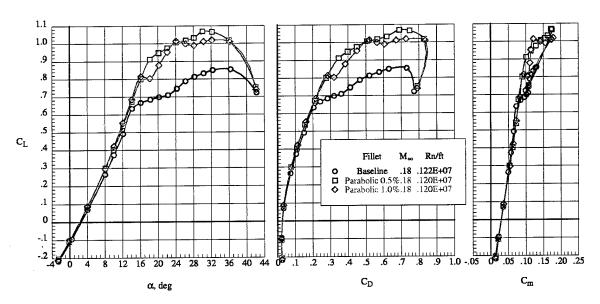


Figure 19- Parabolic Fillet Size Effect on Longitudinal Characteristics, Mach 0.18

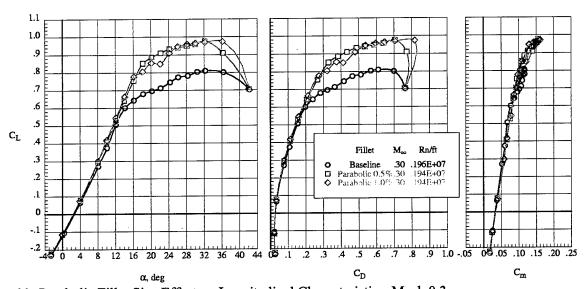


Figure 20- Parabolic Fillet Size Effect on Longitudinal Characteristics, Mach 0.3

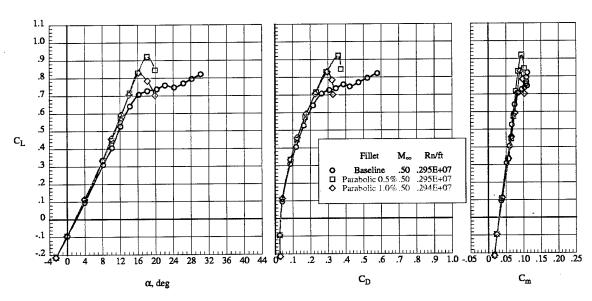
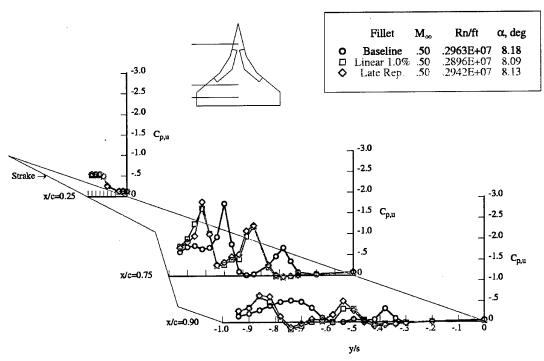


Figure 21- Parabolic Fillet Size Effect on Longitudinal Characteristics, Mach 0.5



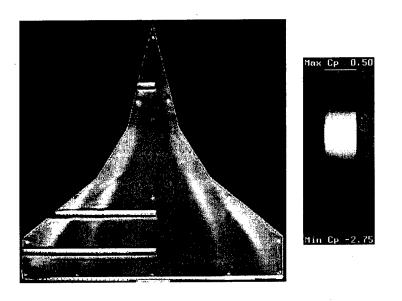


Figure 22- Linear Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA  $8^{\circ}$ 

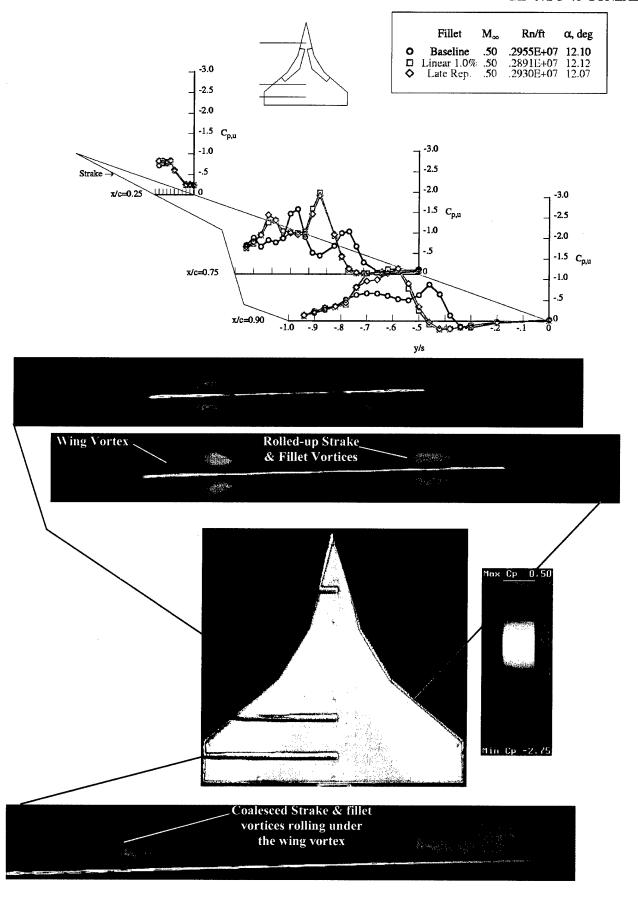


Figure 22 (cont.)- Linear Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 12°

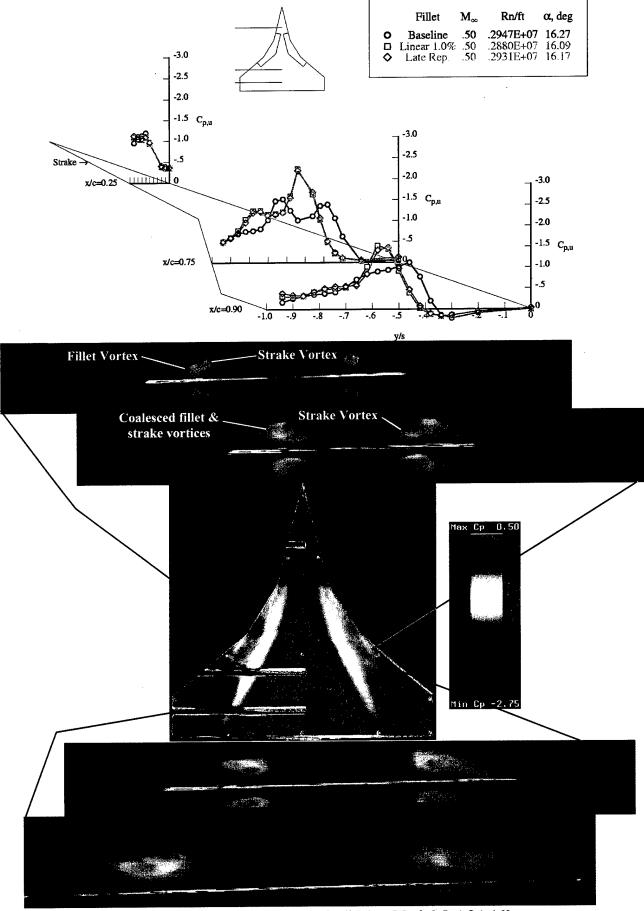


Figure 22 (cont.)- Linear Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 16°

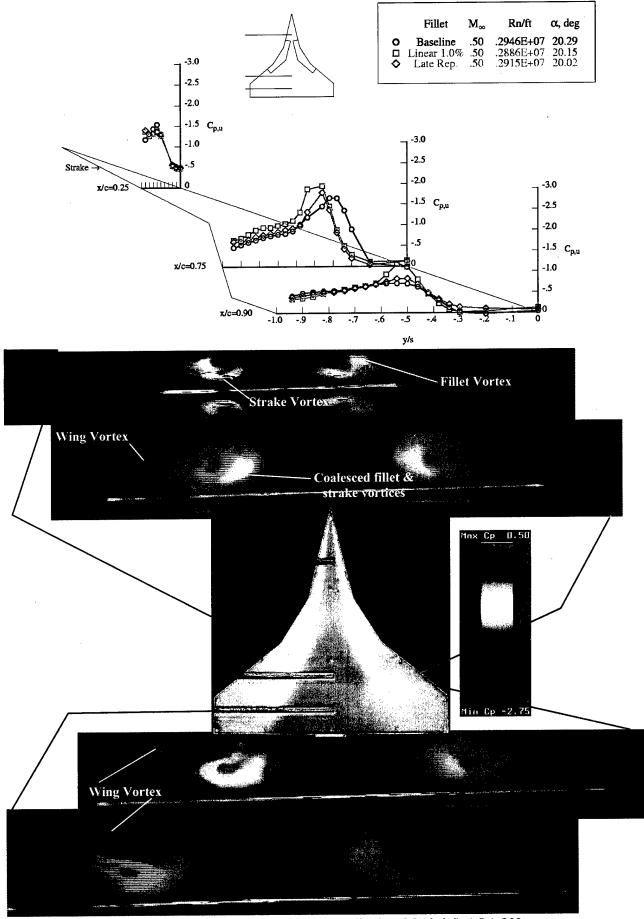
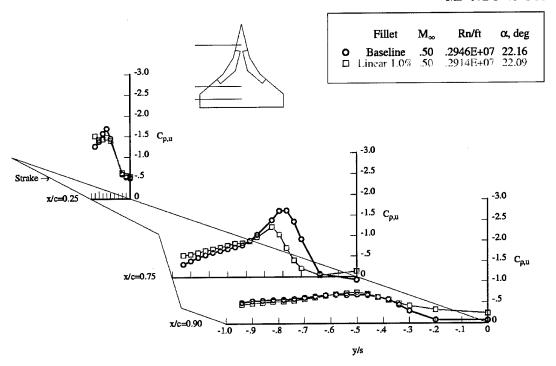


Figure 22 (cont.)- Linear Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 20°



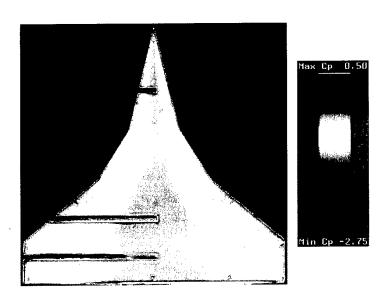
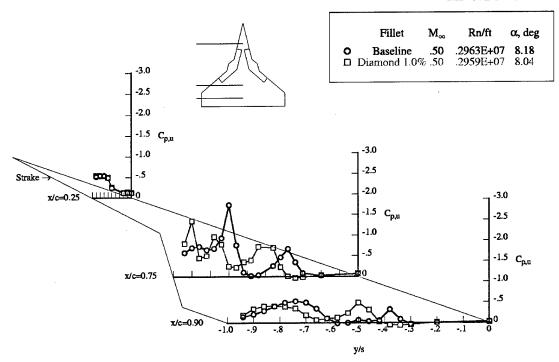


Figure 22 (cont.)- Linear Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA  $22^\circ$ 



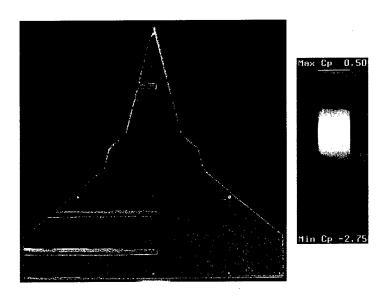
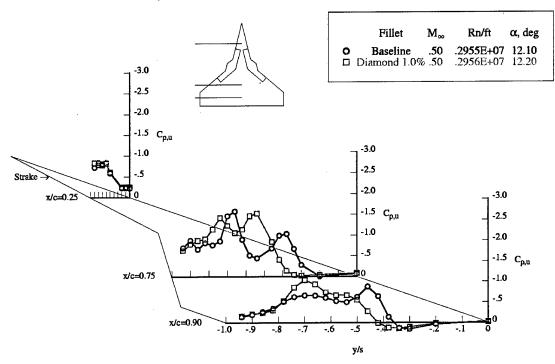


Figure 23- Diamond Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA  $8^{\circ}$ 



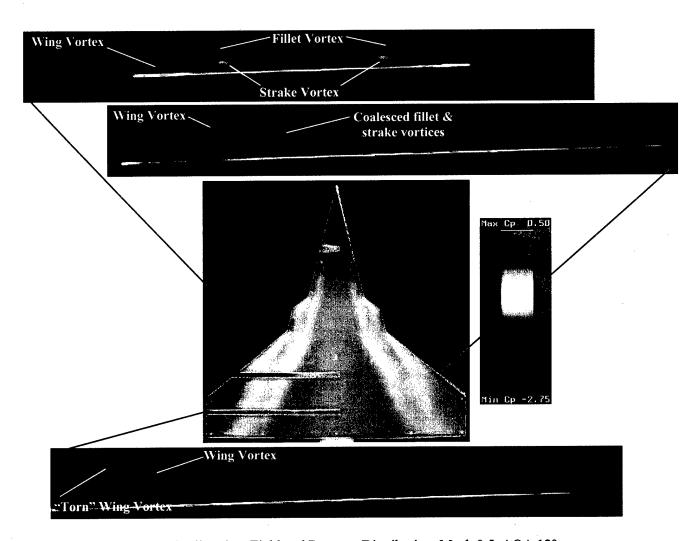


Figure 23 (cont.)- Diamond Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 12°

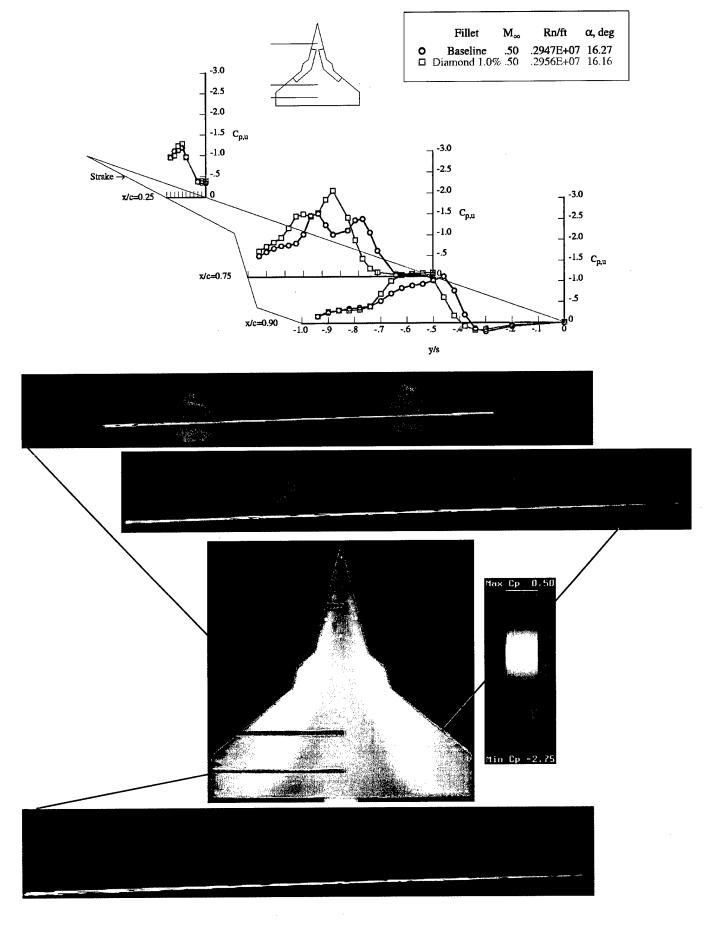


Figure 23 (cont.)- Diamond Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 16°

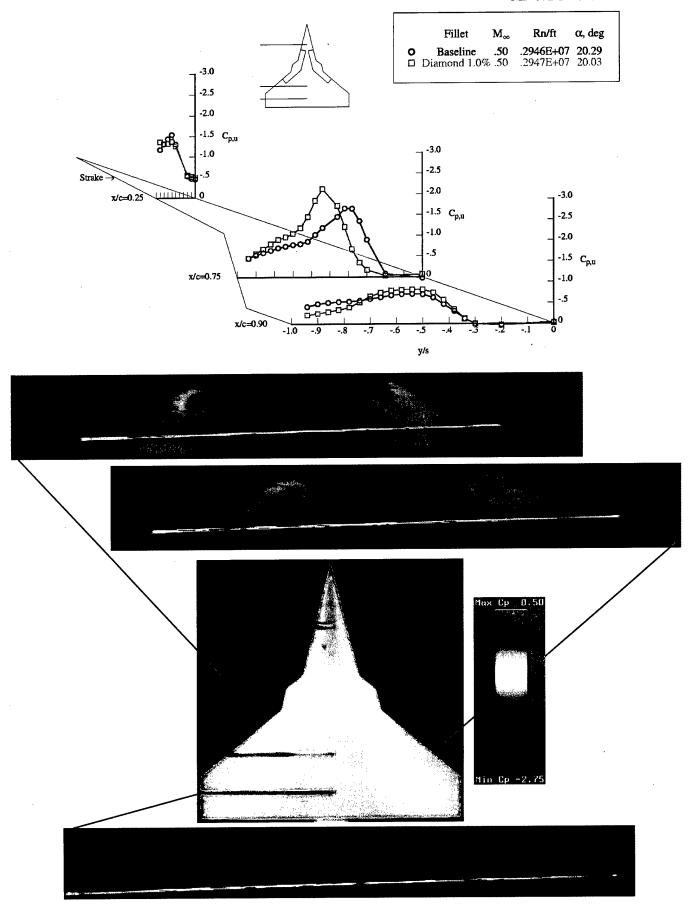


Figure 23 (cont.)- Diamond Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA  $20^\circ$ 

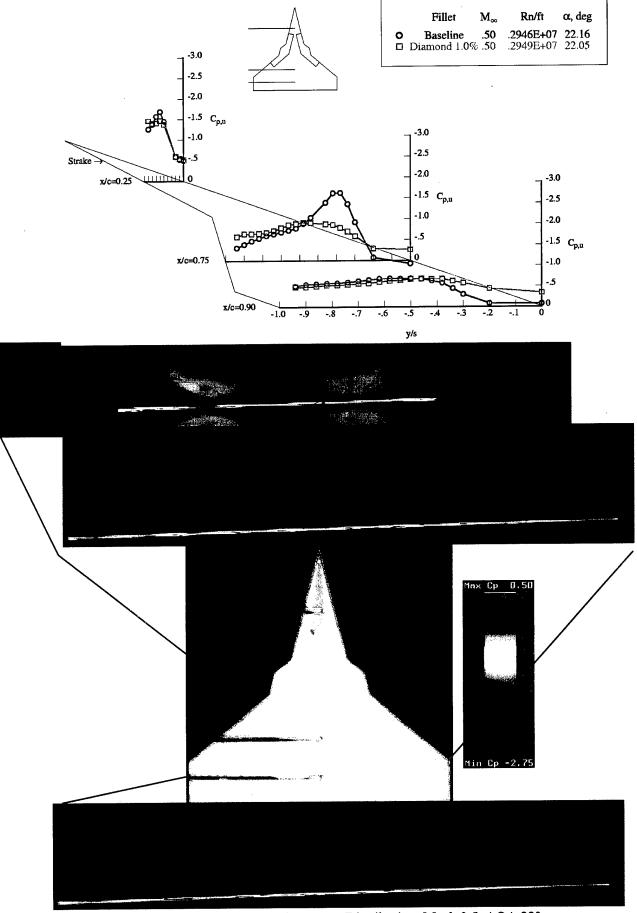
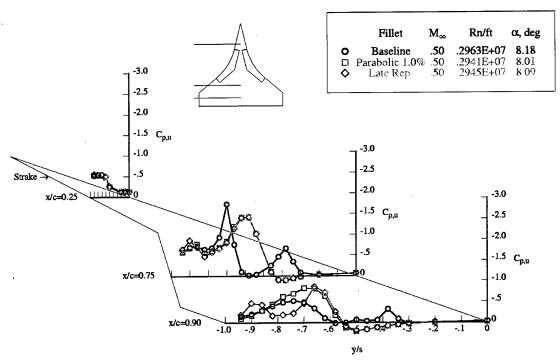


Figure 23 (cont.)- Diamond Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 22°



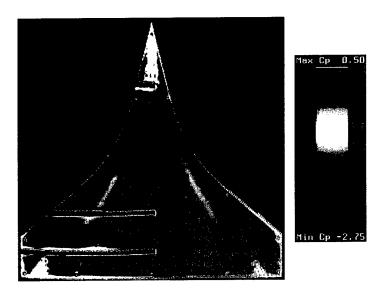


Figure 24- Parabolic Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA  $8^{\circ}$ 

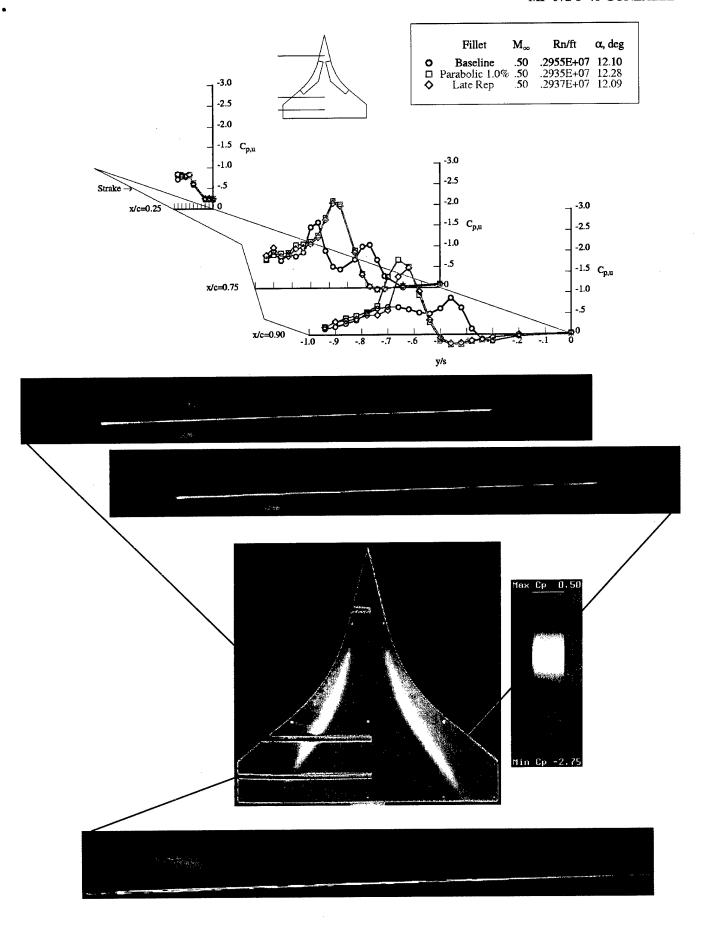


Figure 24 (cont.)- Parabolic Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 12°

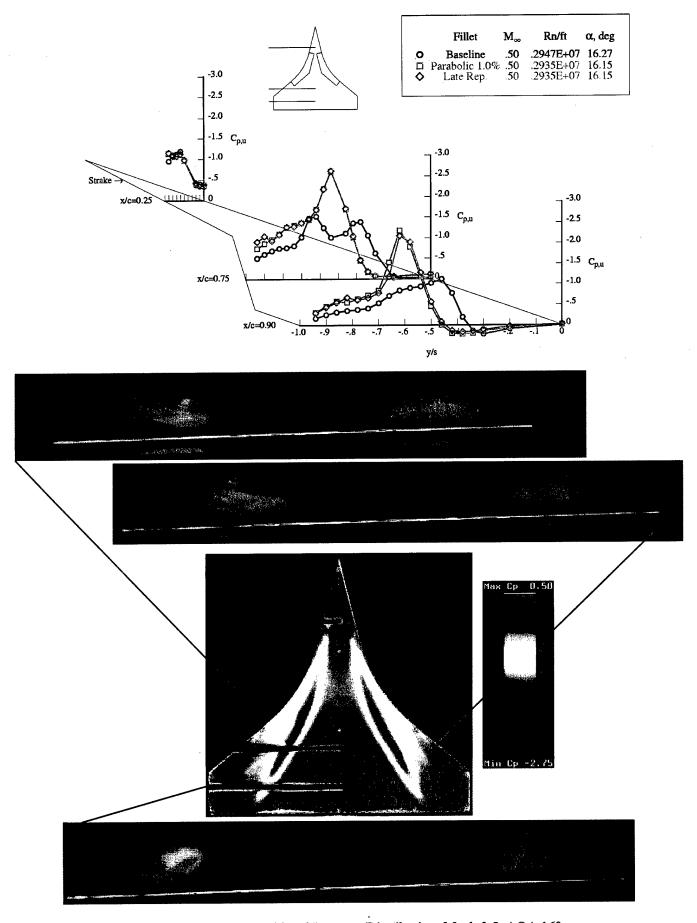


Figure 24 (cont.)- Parabolic Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 16°

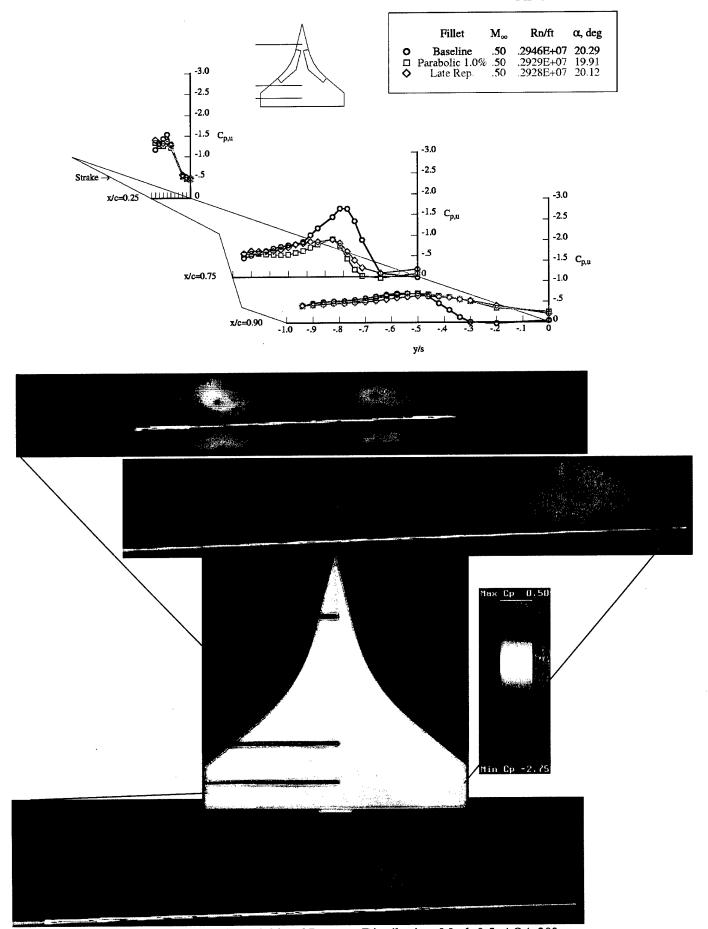


Figure 24 (cont.)- Parabolic Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA 20°

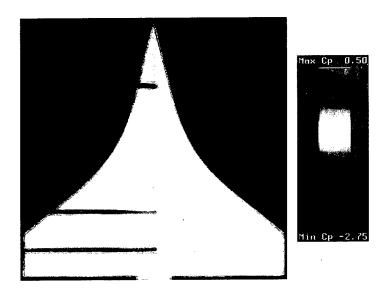


Figure 24 (cont.)- Parabolic Fillet Flow Field and Pressure Distribution, Mach 0.5, AOA  $20^\circ$